

Simulation of positron backscattering on Al, Cu, Ag and Au targets using GEANT4 code

Xin Lai^{1,2}, Xiaopan Jiang¹, Xingzhong Cao^{1ξ}, Xi Zhang^{2}, Zhiming Zhang¹, Xuexiang Cao¹, Gang Xiang², Baoyi Wang¹, Long Wei^{1δ}*

¹Beijing Engineering Research Center of Radiographic Techniques and Equipment, Institute of High Energy Physics, CAS, Beijing 100049, China

²College of Physical Science and Technology, Sichuan University, Chengdu 610064, China

Corresponding authour:

^ξcaoxzh@ihep.ac.cn, ^δweil@ihep.ac.cn, ^{*}xizhang@scu.edu.cn

ABSTRACT

In this paper, backscattering of 3-50 keV positrons on *Al*, *Cu*, *Ag* and *Au* metallic targets has been systematically investigated using GEANT4 code. The dependence of positron backscattering coefficients on atomic number *Z*, target thickness, incident energy and angles has been discussed comprehensively. Besides, positron backscattering spectra for those metallic targets at different discrete scattering angles were also studied to provide theoretical basis of the most appropriate scattering angle selected for simulation parameters and specified applied measurement techniques. The impact of atomic number *Z* of targets on positron backscattering spectra was investigated as well. Simulation results are in reasonable agreement with previous experiment data and theoretical work.

Keywords: GEANT4; Positron; Backscattering coefficients; Backscattering spectra

1. Introduction

The interaction of slow charged particles such as electrons and positrons with solid targets is of prime importance in many applications including electron spectroscopy, electron microlithography, positron annihilation spectroscopy and so on^[1-7]. For instance, slow positron beam techniques have been used as a probe of surfaces or thin films in many ways^[8-12]. Positron backscattering has significant influence on applied measurement results reflecting the information of surface and interfaces of thin films, such as the structure, defects, and electrons distribution of density and momentum. There is a potential that positron beams would be developed as industrial analysis tools utilized to obtain damage distribution information of the surface or near-surface region^[13]. In addition, a comprehensive understanding of the positron collision process in solids underpins and strengthens the description of equivalent electron processes which govern the interpretation of an array of experimental techniques using mono-energetic electrons as detectors^[14]. Backscattered positrons refer to those positrons which, when directed toward the surface of a solid target with some incident energy, are scattered backward with a scattering angle larger than 90° ^[15]. It is well known that the fraction of the primary positrons from surface can be described by the backscattering coefficient, which is usually used as a major index to characterize the absorbed rate of incident particles and could provide basic mechanism to describe the scattering processes accurately^[16]. The dependence of backscattering coefficients for positrons or electrons on target thicknesses, atomic number Z , particle energy and incident angles have been widely investigated^[17-25]. However, most works focused on

one certain factor of those mentioned above, and the corresponding backscattering spectra were rarely reported. In experiments, the energy loss spectra of backscattered positrons could be used to characterize the components, morphology and structure of the material surface.

It is necessary to point out that the direct measurement of positron backscattering coefficients in experiments has some inevitable and uncontrollable factors which affect the experimental results. For instance, the relatively weak intensity of positron beams compared with electron beams and the effect of confining magnetic field on positron beams could affect the results indeed. GEANT4^[26-29] is a very strong toolkit aimed to simulate the passage of particles through condensed matter over a wide energy range. The standard electromagnetic processes related to electron/positron include bremsstrahlung, ionization and δ -ray production, positron annihilation and synchrotron radiation. In the present paper, the simulation results of backscattering coefficients and spectra of positrons impinging on several targets with various atomic number Z and thickness using GEANT4 code are reported. Primary positron energy and incident angles are settled in the range of 3-50 keV and 0° - 80° , respectively. *Al*, *Cu*, *Ag* and *Au* are selected as the targets for the following reasons: atoms are randomly uniformly distributed in the structure established through GEANT4 and the structure properties only depends on the density and its atomic number Z set up in GEANT4. As a result, the models we build is very close to the actual cubic structures of *Al*, *Cu*, *Ag* and *Au* elements.

2. Theories and models

According to the well-known Vicanek and Urbassek theory^[30], the backscattering coefficient could be expressed as:

$$\eta(\theta_i) = \left(1 + a_1 \frac{\mu_0}{\nu^{\frac{1}{2}}} + a_2 \frac{\mu_0^2}{\nu} + a_3 \frac{\mu_0^3}{\nu^{\frac{3}{2}}} + a_4 \frac{\mu_0^4}{\nu^2} \right)^{-1/2} \quad (1)$$

where ν is the mean number of wide-angle collisions defined as:

$$\nu = NR\sigma_{tr} \quad (2)$$

in expression (2), σ_{tr} stands for the transport elastic scattering cross sections:

$$\sigma_{tr} = 2\pi \int_0^\pi (1 - \cos \theta) \frac{d\sigma}{d\Omega} \sin \theta d\theta \quad (3)$$

where θ is the polar-scattering angle^[31].

If the incident mono-energetic positron annihilates with electrons directly during the implantation process, then it is assigned to an “absorbed” positron, if it comes back to the same surface through which it penetrates, then it is assigned to a backscattered positron, and finally, if it crosses through the target, then it is assigned to a transmitted positron. The positrons interact with the atoms by the elastic scattering with the nuclei or the inelastic scattering with the electrons in targets. The scattering processes can be described by the scattering cross sections in the interior of the solid, which could be used to determine the trajectories of the positrons in a GEANT4 simulation^[32]. G4eMultipleScattering are utilized to handle the elastic scattering process. The inelastic cross section is introduced in the present GEANT4 simulation by G4eIonisation to calculate the energy loss. The positron energy decreases gradually in a continuous manner between elastic collisions. The cut value is set to 1 nm and the minimum energy is 20 eV, the amount of positrons set up for each simulation is 10^6 . Fig. 1 shows the model structure used to calculate backscattering spectra. Constant d is set to be 0.7 cm which stand for the distance between the target and the annular detector. We select 1 cm as the thickness of the target which is far greater than its corresponding positron implantation depth. The

target radius is set to be 0.5 cm. the scattering angle θ is adjusted through the control of inside and outside diameters of the annular detector as shown in Table 1.

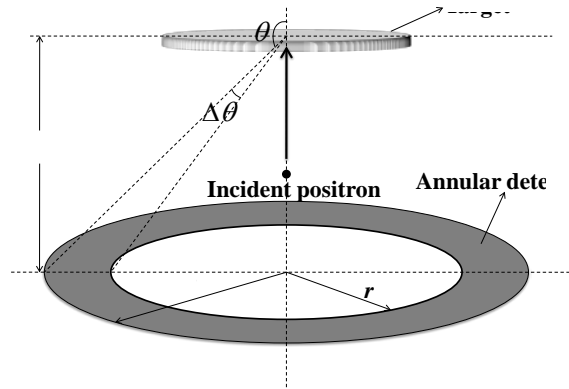


Fig 1. Model constructed for the simulation of positron backscattering spectra.

θ (deg)	Inside (cm)	Outside (cm)
120°-130°	0.8337	1.2124
140°-150°	0.4039	0.5873
160°-170°	0.1232	0.2548

Table 1. The inside and outside diameters of the annular detector for different backscattering angles.

3. Results and discussions

Backscattering coefficients of 3-50 keV positrons impinging on various semi-infinite targets versus incident energy are shown in Fig. 2. Available experimental and theoretical data are presented for comparison, P.G. Coleman et al^[14] have investigated the positron backscattering from elemental solids experimentally. a 72%-efficiency HPGe detector was utilized to record the annihilation positrons through detecting the 0.511 MeV photons produced during annihilation. The amount of backscattered positrons was obtained by removing the annihilation fraction from the total number of positrons. Annihilation count rate C_{Be} for beryllium sample was used to deduce the incident beam intensity I_0 :

$$I_0 = C_{Be} / 0.9625 \quad (4)$$

The backscattering coefficient is then evaluated from:

$$\eta(E) = 1 - 0.9625 C(E) / C(Be) \quad (5)$$

Where $C(E)$ were the annihilation count rates for other samples for each positron energy E selected. As for the theoretical calculation, Chaoui and Bouarissa^[33] have utilized the well-known Vicanek and Urbassek theory mentioned above to calculate the backscattering coefficient of 1-10 keV positrons. The only difference is that the transport cross section σ_{tr} used is expressed as the following series:

$$\sigma_{tr} = 4\pi \sum_{l=0}^{\infty} (l+1) \sin^2(\delta_l - \delta_{l+1}) \quad (6)$$

In expression (6), the phase shifts δ_l are obtained by numerically solving the Schrodinger equation of a positron with an atom bound in a solid.

Our simulation results are reasonably agreeable with them in similar variation tendency, the deviations between simulation results and the experimental data may be introduced by experimental errors or GEANT4 code itself. Actually, the detailed crystal structure of the host material has not been perfectly incorporated in GEANT4 code. Although the model structures (*Al*, *Cu*, *Ag* and *Au*) established are very close to their original cubic close packed structures, lack of the crystal effect still has an impact on the results to some extent.

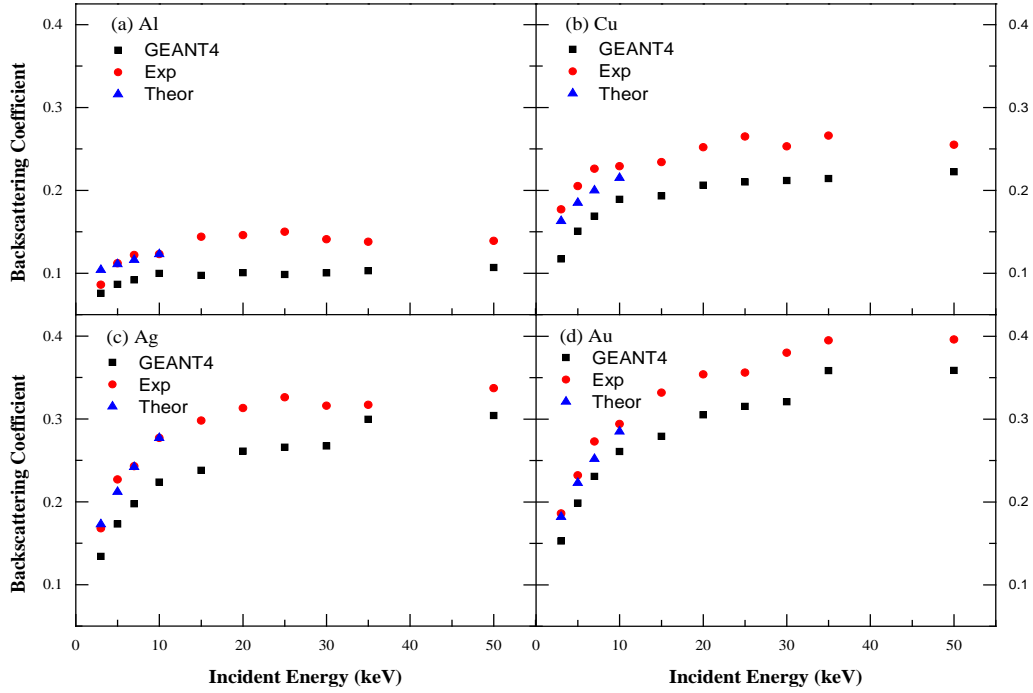


Fig 2. Backscattering coefficients of 3-50 keV positrons on various targets. (a) *Al*, (b) *Cu*, (c) *Ag* and (d) *Au*. experimental (Exp)^[14], and theoretical (Theor)^[33] results are presented for comparison.

Table 2. GEANT4, experimental (Exp)^[14] and theoretical (Theor)^[33] results for positron backscattering coefficients from *Al*, *Cu*, *Ag* and *Au*.

<i>Z</i>	<i>E</i> (keV)	3	5	7	10	15	20	25	30	35	50
13	GEANT4	0.076	0.086	0.092	0.099	0.097	0.100	0.098	0.100	0.103	0.107
	Theor	0.104	0.111	0.116	0.123	-	-	-	-	-	-
	Exp	0.086	0.112	0.122	0.123	0.144	0.146	0.15	0.141	0.138	0.139
29	GEANT4	0.117	0.151	0.169	0.189	0.193	0.206	0.210	0.212	0.214	0.222
	Theor	0.163	0.185	0.2	0.215	-	-	-	-	-	-
	Exp	0.177	0.205	0.226	0.229	0.234	0.252	0.265	0.253	0.266	0.255
47	GEANT4	0.134	0.173	0.198	0.224	0.238	0.261	0.266	0.267	0.299	0.304
	Theor	0.173	0.212	0.242	0.277	-	-	-	-	-	-
	Exp	0.168	0.227	0.243	0.277	0.298	0.313	0.326	0.316	0.317	0.337
79	GEANT4	0.153	0.199	0.231	0.261	0.279	0.305	0.315	0.321	0.359	0.359
	Theor	0.182	0.223	0.252	0.285	-	-	-	-	-	-
	Exp	0.186	0.232	0.273	0.294	0.332	0.354	0.356	0.38	0.395	0.396

Table 3. Other MC simulation results^[33-34] of backscattering coefficients of 1-10 keV positrons on *Al*, *Cu*, *Ag* and *Au*.

<i>Z</i>	<i>E</i> (keV)	1	2	3	4	5	6	7	8	9	10
13	MC[33]	0.086	0.098	0.104	0.107	0.111	0.113	0.116	0.118	0.120	0.123
	MC[34]	0.095	0.098	0.098	0.096	0.099	0.099	0.100	0.104	0.105	0.109
29	MC[33]	0.117	0.146	0.163	0.175	0.185	0.193	0.200	0.205	0.210	0.215
	MC[34]	0.148	0.169	0.178	0.188	0.194	0.199	0.208	0.210	0.214	0.218
47	MC[33]	0.109	0.146	0.173	0.194	0.212	0.228	0.242	0.254	0.266	0.277
	MC[34]	0.131	0.154	0.177	0.198	0.216	0.230	0.241	0.255	0.264	0.278
79	MC[33]	0.112	0.153	0.182	0.205	0.223	0.239	0.252	0.264	0.275	0.285
	MC[34]	0.128	0.161	0.189	0.207	0.226	0.240	0.251	0.263	0.273	0.276

As shown in Figure. 2 and Table. 2, positron backscattering coefficients present a monotonic increase as the target atomic number Z and incident energy increase, which could be explained by the fact that heavy atoms have a larger probability of scattering into large angles than the light atoms in elastic collisions, namely, a larger scattering cross section. It is worth noting that the backscattering coefficient increases gradually until the incident energy increases to a certain value (which depends on the target atomic number Z) and then becomes relatively constant for each metal. This suggests that the increase of positron incident energy may lead to the increase of scattering cross section in a particular energy range.

As shown in Table. 3, we also presented other MC simulation results for positron ($E=1-10$ keV) backscattering coefficients from *Al*, *Cu*, *Ag* and *Au*. It is found that GEANT4 results is more close to the experiment compared with other MC simulation results when the backscattering coefficient is relatively small (~ 0.100). This suggests GEANT4 simulation might be more reliable when the positron has a very low incident energy. Actually, the scattering model we selected leads to the difference of results obtained by various MC strategies. For example, ref-34 adopts the optimized Wentzel (OW2) model^[35], which is different with the our G4eMultipleScattering and G4eIonisation model. The elastic cross sections used in the optimized Wentzel (OW2) model are obtained utilizing the Dirac PW analysis (ELESPE code)^[36] and the inelastic process is introduced through using a simple approximation to calculate the energy loss of the positron. Actually, ref-12 has also investigated the positron ($E=1-5$

keV) backscattering coefficients from semi-infinite *Al*, *Ag* and *Au* through MC simulation, the results are in reasonable agreement with the present GEANT4 work.

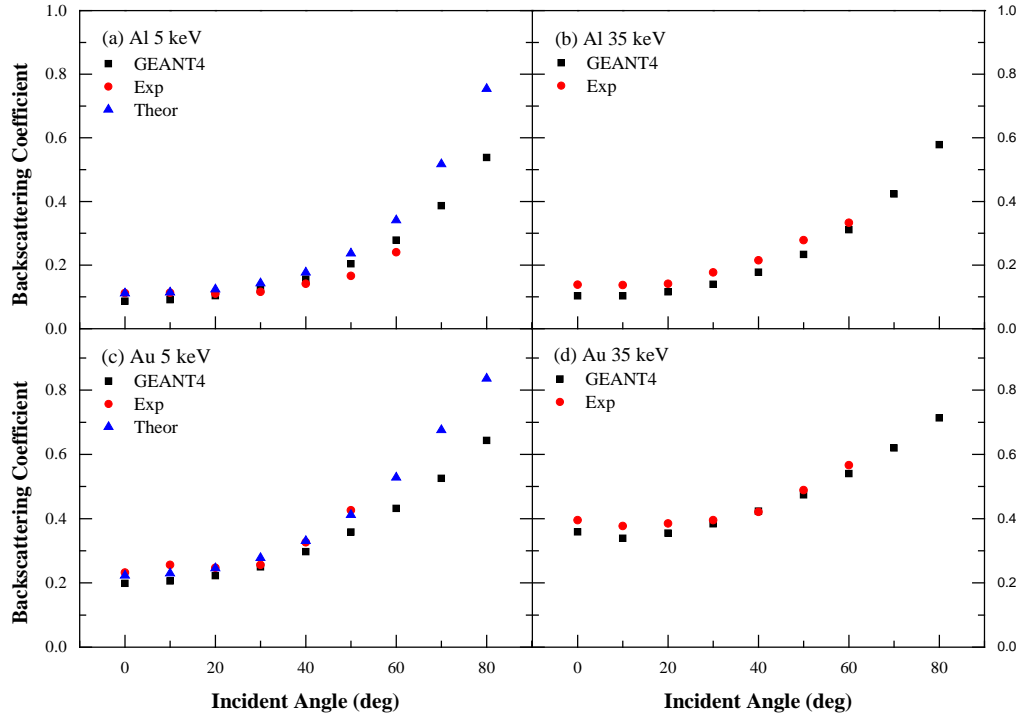


Fig 3. Backscattering coefficients of positrons impinging on *Al* and *Au* targets versus incident angles. (a) *Al*, 5.0 keV (b) *Al*, 35.0 keV (c) *Au*, 5.0 keV and (d) *Au*, 35.0 keV. experimental (Exp)^[14], and theoretical (Theor)^[33] results are presented for comparison.

Table 4. GEANT4, experimental (Exp)^[14], theoretical (Theor)^[33] and MC simulation (MC)^[34] results for 5 keV positron backscattering coefficients from *Al* and *Au* versus incident angles.

Z	θ (deg)	0	10	20	30	40	50	60	70	80
13	GEANT4	0.086	0.091	0.104	0.123	0.155	0.204	0.280	0.386	0.538
	Theor	0.111	0.114	0.124	0.143	0.177	0.237	0.341	0.517	0.754
	Exp	0.112	0.113	0.11	0.116	0.141	0.166	0.24	-	-
	MC	0.126	0.134	0.143	0.167	0.203	0.254	0.334	0.443	0.568
79	GEANT4	0.199	0.206	0.223	0.249	0.297	0.358	0.432	0.525	0.643
	Theor	0.223	0.23	0.246	0.278	0.331	0.412	0.528	0.676	0.836
	Exp	0.232	0.256	0.247	0.256	0.326	0.426	-	-	-
	MC	0.290	0.295	0.301	0.336	0.369	0.422	0.478	0.552	0.673

Figure. 3 and Table. 4 illustrates the dependence of positron backscattering

coefficients on the incident angle θ ($0^\circ \leq \theta \leq 80^\circ$ relative to the surface normal of the target species) for semi-infinite *Al* (light element) and *Au* (heavy element). 5.0 and 35.0 keV are selected as the typical incident energy for each target. The available experimental, MC simulation and theoretical results are also presented for comparison. The GEANT4 results indicate that backscattering coefficients of positrons impinging on both *Al* and *Au* targets increase with the incident angle θ (independently of incident energy). It is worth noting that backscattering coefficients increase slowly below 20° , but vary rapidly when incident angles becomes larger. The discrepancies of the value of positron backscattering coefficients in these four situations are smaller and smaller. It is found that the backscattering coefficient of oblique incidence presents a significant increase compared with that of normal incidence, especially when the incident angle becomes larger. namely, it is easier for positrons to scatter through larger angle than 90° compared to the case of normal incidence. This could be attributed to the fact that positrons have encountered several small angle scatterings and lost more energy during the longer traversed path in targets and the small total elastic scattering cross section^[22]. As a result, the positron penetration depth becomes shallow, namely, the corresponding escape distance for positrons decreases, which causes positrons to emit out of the target surface more easily and leads to the increase of the positron backscattering coefficient at non-normal incidence. The results demonstrate that the positron backscattering coefficient is related to the incident angles and selecting an appropriate incident angle for measurement techniques could improve the effective counts and reduce the noise of background.

For 35.0 keV positrons with normal incidence, the variation of the backscattering coefficient ratio with the thickness of targets is shown in Fig. 4. Where, $Bsc(t)$ is referred to a given thickness of targets and $Bsc(si)$ is referred to the semi-infinite

targets. The results show that for a certain material, the positron backscattering coefficient increases gradually until the target thickness increases to a critical value (which depends on the target material) and then becomes almost constant^[15,21,37-38]. This could be explained by the simple fact that if the target is a bulk material, then the number of transmitted positrons would be zero and the fraction of backscattered positrons could reach its maximum value (backscattering coefficients of semi-infinite targets). It could be noticed that the critical thickness decreases with the increase of atomic number Z , which indicates that the impact of target surface on the backscattering coefficient is also related to the atomic number Z . Actually, the mean positron implantation depths for semi-infinite materials are customarily expressed as: $\bar{z} = \frac{A}{\rho} E^n$, where E is the positron energy, ρ is the material density, and A and n are energy-independent constants for a particular material^[39-40]. It could be speculated that the positron implantation depth is deeper for lighter atoms and hence a larger transmission probability considering the same target thickness. As a result, the possibility of emitting out of the target surface for incident positrons becomes smaller.

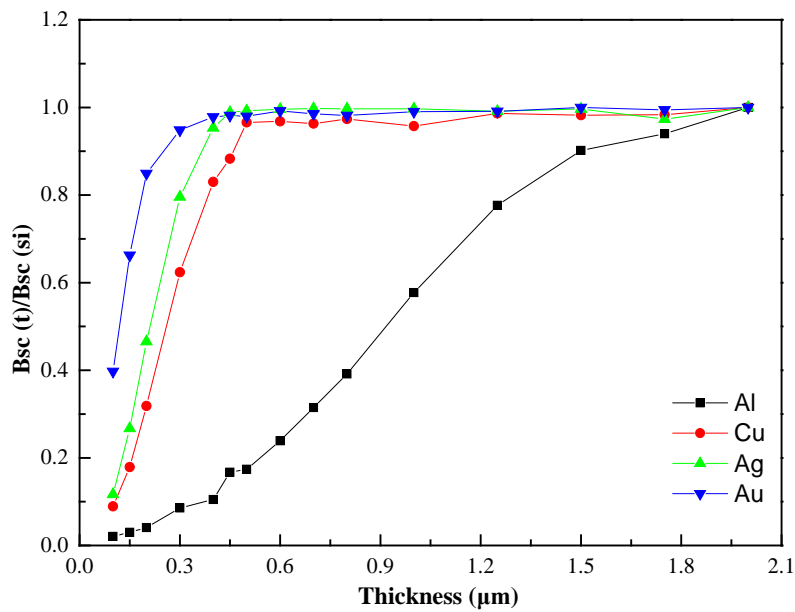


Fig 4. The ratio of positron backscattering coefficient for the target thickness to that of the semi-infinite target, $Bsc(t)/Bsc(si)$ versus the target thickness for 35 keV positrons.

Positron backscattering spectra for Au targets are simulated at all angles in the range $100^\circ \leq \theta \leq 180^\circ$ at 10° intervals. It is relatively difficult to extract data for scattering angles less than 120° or more than 170° . Fig. 5 and Fig. 6 present the backscattering spectra and relative (energy-integrated) backscattering yields of positrons impinging on Au targets with normal incidence, respectively. It is found that that the most probable energy of backscattered positrons changes very little at different scattering angles. Besides, scattering angles in the range 120° - 140° have the maximal backscattering yields, backscattering yields decrease when scattering angles are less than 110° or more than 150° . This could provide theoretical support for parameters setting and specified applied measurement techniques.

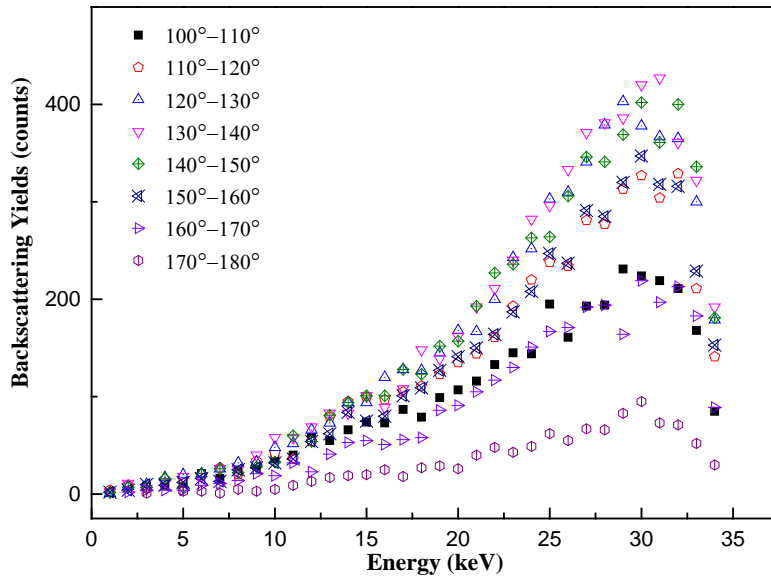


Fig 5. Backscattering spectra of 35 keV positrons impinging on semi-infinite Au targets at different discrete scattering angles.

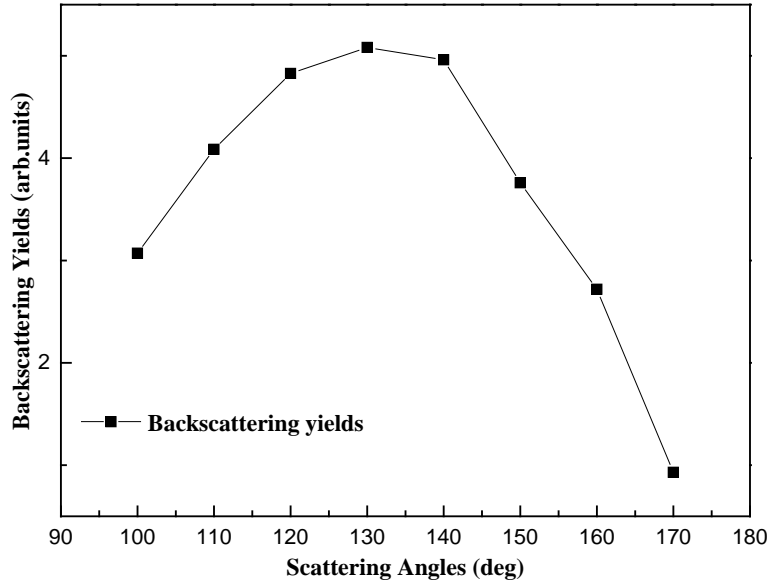


Fig 6. Relative backscattered yields (integrated over energy) of 35 keV positrons impinging on semi-infinite *Au* targets.

In order to investigate the impact of atomic number Z and scattering angles on the energy distribution, energy spectra of backscattered positrons with a normal incidence impinging on semi-infinite *Al*, *Cu*, *Ag* and *Au* targets are simulated. On the basis of the parameters setting mentioned above, we select three typical discrete scattering angles 120° - 130° , 140° - 150° and 160° - 170° . As shown in Fig. 7, It is found that the most probable energy of backscattered positrons changes very little at different scattering angles considering the same target species, but varies a lot for different target material considering the same scattering angle. This suggests that the influence of scattering angles on the energy distribution of backscattered positrons is less pronounced than that of atomic number Z . The peak of positron backscattering yields increases with target atomic number Z considering the same scattering angle, which is consistent with the fact that the increase of target atomic mass leads to less energy loss and larger scattering cross section. Fig. 8 gives the relative (energy-integrated) backscattered yields versus the atomic number Z . The results show that the backscattering yields vary very little from 120° - 130° to 140° - 150° and the decrease at

160°-170° is significant compared to the former, which agrees with the results shown in Fig. 6 very well.

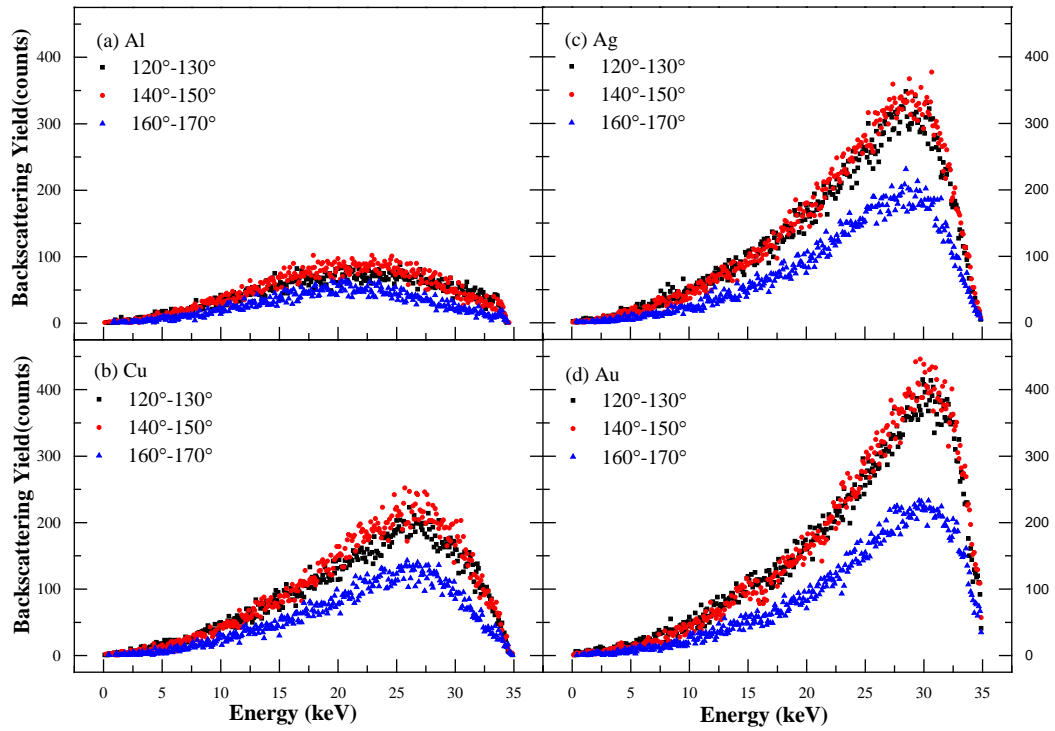


Fig 7. Backscattering spectra of 35 keV positrons on various semi-infinite targets at three discrete scattering angles 120°-130°, 140°-150° and 160°-170°. 10^6 incident positrons are used in each simulation run. (a) *Al*, (b) *Cu*, (c) *Ag* and (d) *Au*.

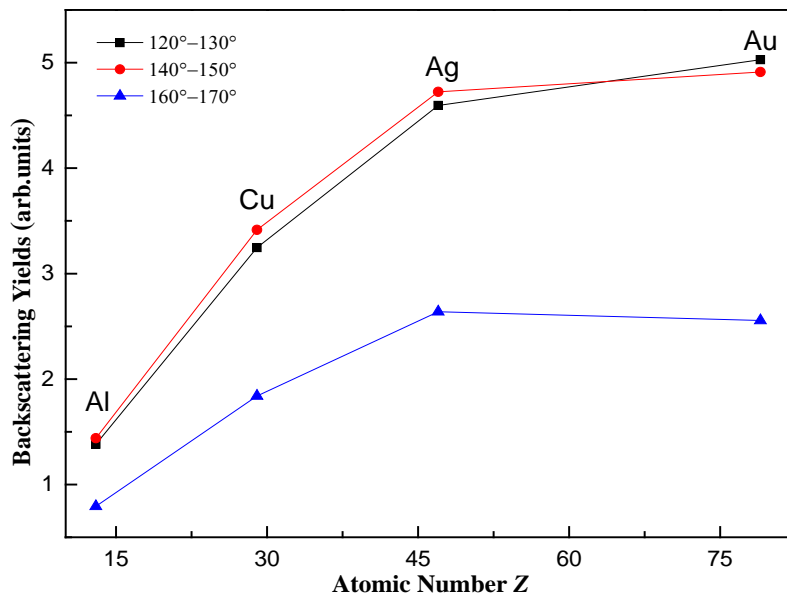


Fig 8. Relative backscattered yields (integrated over energy) of 35 keV positrons on semi-infinite *Al*, *Cu*, *Ag* and *Au* targets at three discrete scattering angles 120°-130°, 140°-150° and 160°-170°. The solid lines are drawn to guide the eye.

4. Conclusions

In the present paper, backscattering coefficients and spectra of 3–50 keV positrons on a variety of targets (*Al*, *Cu*, *Ag* and *Au*) with various thickness for different incident angles are simulated using GEANT4 code. The results show reasonable agreement with previous experiment data and theoretical calculation. It is found that positron backscattering coefficient presents a monotonic increase with both incident angles and atomic number *Z* of targets. Considering a certain primary positron energy, the critical value of the target thickness at which backscattering coefficients become constant depends on the target material. The investigation of backscattering spectra indicates that atomic number *Z* has a greater impact on the energy distribution of backscattered positrons compared to scattering angles. Our results may be useful for the application of slow positrons as a probe for obtaining information on the surface or near-surface region.

Acknowledgments

This work is supported by the National Natural Science Foundation (Grant Nos. 11475197, 11475193, 11175191).

References

1. P.J. Schultz, K.G. Lynn, *Rev. Mod. Phys.* **1988**, 60, 701.
2. P.G. Coleman, *Appl. Surf. Sci.* **2002**, 194, 264.

3. C.J. Powell, A. Jablonski, S. Tanuma, *J. Electron Spectrosc. Relat. Phenom.* 1994, **68**, 605.
4. I. Adesida, R. Shimizu, T.E. Everhart, *J. Appl. Phys.* **1980**, 51, 5962.
5. W.S.M. Werner, *Surf. Interface Anal.* **2001**, 31, 141.
6. J. Rundgren, *Phys. Rev. B.* **1999**, 59, 5106.
7. G. Gergely, *Prog. Surf. Sci.* **2002**, 71, 31.
8. G.R. Massoumi, N. Hozhabri, K.O. Jensen, W.N. Lennard, M.S. Lorenzo, P.J. Schultz, A.B. Walker, *Phys. Rev. Lett.* **1992**, 68, 3873.
9. M. Dapor, *J. Appl. Phys.* **1995**, 77, 2840.
10. N. Bouarissa, A.B. Walker, H. Aourag, *J. Appl. Phys.* **1998**, 83, 3643.
11. Z. Chaoui, N. Bouarissa, *J. Appl. Phys.* **2004**, 96, 807.
12. N. Bouarissaa, M.S. Al-Assiri, *J. Electron. Spectrosc. Relat. Phenom.* **2013**, 191, 11.
13. A. Zecca, *Appl. Surf. Sci.* **2002**, 194, 4.
14. P.G. Coleman, L. Albrecht, K.O. Jensen, A.B. Walker, *J. Phys.: Condens. Matter.* **1992**, 4, 10311.
15. A. Bentabet, N. Bouarissa, *Appl. Surf. Sci.* **2007**, 253, 8725.
16. J. Makinen, S. Palko, J. Martikainen, P. Hautjarvi, *J. Phys.: Condens. Matter.* **1992**, 4, L503.
17. V.J. Ghosh, G.C. Aers, *Phys. Rev. B.* **1995**, 51, 45.
18. M. Dapor, *J. Appl. Phys.* **1996**, 79, 8406.
19. G.R. Massoumi, W.N. Lennard, P.J. Schultz, A.B. Walker, K.O. Jensen, *Phys. Rev. B.* **1993**, 47, 11007.
20. K.O. Jensen, A.B. Walker, *Surf. Sci.* **1993**, 292, 83.
21. M. Dapor, *Phys. Rev. B.* **1992**, 46, 618.

22. N. Bouarissa, B. Deghfel, A. Bentabet, *Europ. Phys. J. AP.* 2002, 19, 89.
23. G.R. Massoumi, N. Hozhabri, W.N. Lennard, P.J. Schultz, *Phys. Rev. B.* 1991, 44, 3486.
24. A.P. Knights, P.G. Coleman, *J. Phys.: Condens. Matter.* **1995**, **7**, 3485.
25. J.A. Baker, P.G. Coleman, *J. Phys. C: Solid State Phys.* **1988**, 21, L875.
26. S. Agostinelli, J. Allison, K. Amako, et al, *Nucl. Instrum. Methods Phys. Res., Sect. A.* **2003**, 506, 250.
27. J. Dryzek, P. Horodek, *Nucl. Instrum. Methods Phys. Res., Sect. B.* **2008**, 266, 4000.
28. S. Agostinelli et al., *Nucl. Instr. and Meth. A*, **2003**, 506, 250.
29. V. N. Ivanchenko, *Nucl. Instr. and Meth. A*, **2003**, 502, 666.
30. M. Vicanek, H.M. Urbassek, *Phys. Rev. B.* **1991**, 44, 7234.
31. A. Jablonski, *Phys. Rev. B.* **1998**, 58, 24.
32. S. Valkealahti, R.M. Nieminen, *Appl. Phys. A.* **1983**, 32, 95.
33. Z. Chaoui, N. Bouarissa, *J. Phys.: Condens. Matter.* **2004**, 16, 799.
34. Z. E. A. Chaoui, *J. Phys.: Condens. Matter.* **2006**, 18, 10303
35. Z. Chaoui, *Appl. Phys. Lett.* **2006**, 88, 024105
36. F. Salvat, A. Jablonski, C. J. Powell, *Comput. Phys. Commun*, **2005**, 165, 157
37. B. Deghfel, A. Bentabet, N. Bouarissa, *Phys. Status Solidi (b).* **2003**, 238, 136.
38. M. Dapor, *Eur. Phys. J. Appl. Phys.* **2002**, 18, 155.
39. V. J. Ghosh, M. McKeown, D.O. Welch, KG. Lynn, *Nucl. Instrum. Methods Phys. Res., Sect. B.* **1994**, 90, 442.
40. N. Bouarissa, A. B. Walker, *Int. J. Mod. Phys. B*, **2000**, 14, 1603